

80-dB Microwave Noise from an Avalanche Transistor Circuit

Gytis Mykolaitis¹, Saulius Kersulis², Skaidra Bumeliene², Arunas Tamasevicius²

¹*Department of Physics, Vilnius Gediminas Technical University, Sauletekio Al. 11, LT-10223 Vilnius, Lithuania*

²*Institute of Semiconductor Physics, Center for Physical Sciences and Technology, A. Gostauto St. 11, LT-01108 Vilnius, Lithuania*
 gytis.mykolaitis@vgtu.lt

Abstract—Extremely high ‘excess noise ratio’ of 80 dB is observed from an avalanche transistor circuit, operating in a random pulse mode. Broadband noise spectrum measured from 30 MHz to 1 GHz exhibits good flatness with the nonuniformity of only ≈ 1 dB. Experiments have been performed with the silicon bipolar junction microwave transistors. An analog circuit model is proposed and investigated.

Index Terms—Bipolar transistor circuits, microwave transistors, avalanche breakdown, noise generators.

I. INTRODUCTION

Electronic circuits, employing bipolar junction transistors biased into the avalanche breakdown region are known for many years [1], [2]. A variety of circuit modifications with a number of applications to switching devices, nanosecond and subnanosecond pulse generators have been developed. Commonly the pulse generators, when supplied with voltage V_0 noticeably above the threshold voltage V_{th} of the collector-base junction ($V_0 > V_{th}$), produce periodic pulses [3]–[7]. In this paper, we describe an avalanche transistor circuit operating in a random pulse mode, which is observed at supply voltages very close to the threshold, $V_0 \approx V_{th}$.

II. EXPERIMENTAL RESULTS

The circuit diagram is sketched in Fig. 1. The transistor Q is a silicon bipolar microwave BFG520 type ($f_T = 9$ GHz, $V_{th} \approx 26$ V) device. The values of the passive elements are the following: $C_1 = 1$ pF, $R_1 = 24$ k Ω , $R_2 = 1$ k Ω , $R_3 = 50$ Ω .

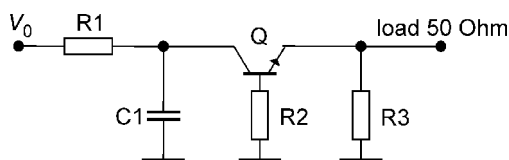


Fig. 1. Avalanche transistor circuit.

Pulse trains, recorded by means of a digital oscilloscope Tektronix DPO71604C (bandwidth 16 GHz), are presented in Fig. 2 and Fig. 3, for random and periodic mode of operation, respectively (note different time scales in Fig. 2 and Fig. 3). An individual pulse is shown in Fig. 4. Its

amplitude and shape do not depend on the applied voltage; they are the same in the random and the periodic modes.

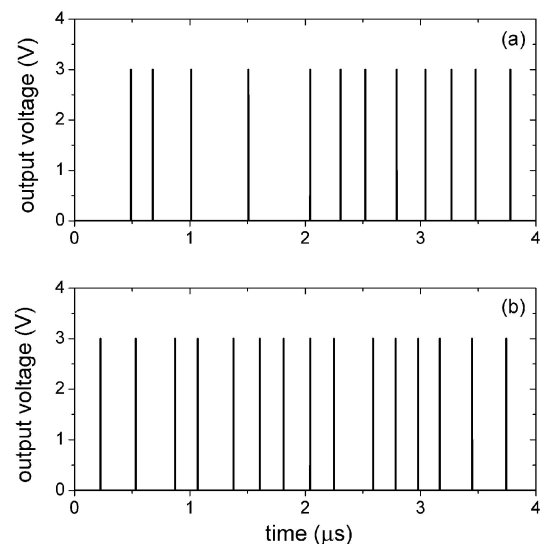


Fig. 2. Two snapshots (a) and (b) of random pulses. $V_0 = 26$ V.

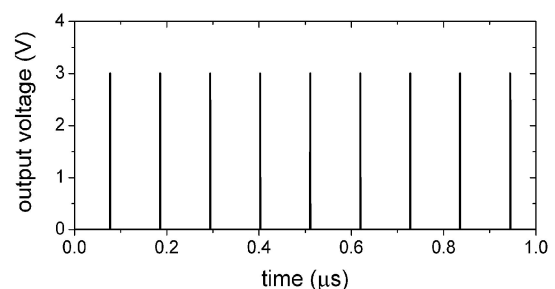


Fig. 3. Periodic pulse train. $V_0 = 28.5$ V.

For the pulse trains examined above, corresponding noise level commonly can be characterized by the power spectral density $k_B T_n$ (in W/Hz), where k_B is the Boltzmann constant, T_n is the effective noise temperature (in K). The output level in the specifications of noise devices is conventionally characterized by the ‘excess noise ratio’ (ENR) either in a linear or in a logarithmic scale, respectively

$$ENR = \frac{T_n - T_0}{T_0}, \quad (1)$$

$$ENR = 10 \log_{10} \frac{T_n - T_0}{T_0}, \quad [\text{dB}]. \quad (2)$$

where T_0 is the standard temperature (290 K).

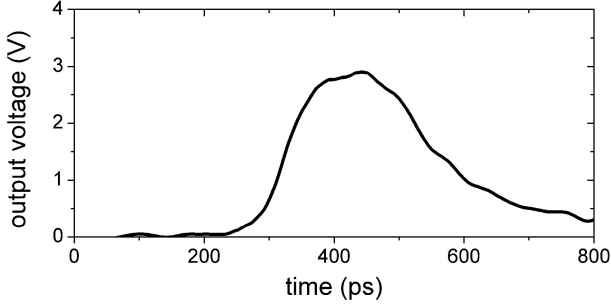


Fig. 4. Experimentally recorded individual pulse.

The spectral distribution of the $ENR(\tilde{S})$ can be estimated from the shape of an individual pulse (Fig. 4), which we approximate by the Gaussian function

$$V(t) = V_m \exp \left[-\frac{1}{2} \left(\frac{t - t_0}{\ddagger} \right)^2 \right], \quad (3)$$

where $V_m = 3$ V, $t_0 = 430$ ps, and $\ddagger = 100$ ps. Then the envelope of the ENR spectrum is given by

$$ENR(\tilde{S}) = ENR(0) \exp \left(-\frac{\tilde{S} \ddagger}{2} \right), \quad (4)$$

where $\tilde{S} = 2\pi f$. Based on the above analytical expression, the spectral density at $f = 1$ GHz decreases to the level of $0.8 \times ENR(0)$.

The experimental power spectrum is presented in Fig. 5.

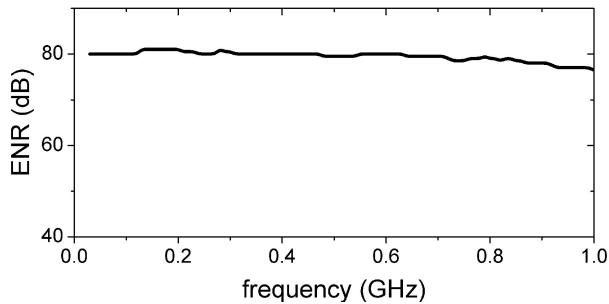


Fig. 5. Experimental power spectrum of random pulses. $V_0 = 26.2$ V. Spectral resolution 120 kHz.

TABLE I. NOISE PARAMETERS FOR DIFFERENT SOURCES.

Noise source	Ref.	T_n (K)	ENR (dB)
Gas discharge tube	[8]	$2.5 \times 10^3 - 3 \times 10^4$	8.8 – 20.1
Avalanche diode	[9]	$1.5 \times 10^5 - 1 \times 10^6$	26 – 35
Microplasma diode	[10]	1×10^8	55
Avalanche transistor	Fig. 5	3×10^{10}	80
Chaotic oscillator	[11], [12]	3×10^{11}	90

The power spectrum (Fig. 5) demonstrates extremely high, 80 dB spectral density and rather flat spectral distribution (± 1 dB). The decrease to the level of $0.5 \times$

$ENR(0)$, i.e. by -3 dB at 1 GHz is in a reasonable agreement with the estimation from (3). Noise level (T_n and ENR) for the avalanche transistor circuit is compared with some other noise sources in Table I.

The microplasma diode in Table I refers to GaAs avalanche diode, biased below the breakdown voltage and periodically driven by external microwave voltage at 1 GHz [10]. The chaotic oscillator [11], [12] is the two-stage Colpitts circuit with the same microwave silicon BFG520 type transistors, operating in the 300 MHz–1000 MHz range.

III. ANALOG MODEL

To demonstrate the mechanism behind the observed random pulses in the avalanche transistor circuit, we propose a phenomenological analog model, which is presented in Fig. 6 as an electronic circuit. The R1–C1 subcircuit plays the same role of charge accumulating unit as in the experimental circuit (Fig. 1). The operational amplifier OA stage is a common buffer used to minimize leakage from the capacitor C1. The comparator unit COMP with a positive feedback via the resistors R2 and R3 implements a hysteretic Schmitt trigger [13]. The ON and the OFF threshold input voltages of the Schmitt trigger are given by:

$$V_{on} = V_h \frac{R_2}{R_2 + R_3} + (E + V_{noise}) \frac{R_3}{R_2 + R_3}, \quad (5)$$

$$V_{off} = V_l \frac{R_2}{R_2 + R_3} + (E + V_{noise}) \frac{R_3}{R_2 + R_3}. \quad (6)$$

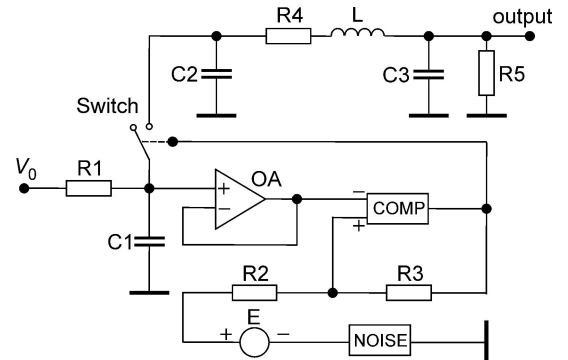


Fig. 6. Analog low frequency circuit, modelling the microwave avalanche transistor circuit in Fig. 1.

In (5), (6) the V_h and the V_l are the high level and the low level output voltage of the comparator, respectively (specifically, $V_h = 13$ V, $V_l = 0.2$ V), the E is the dc bias voltage from the voltage source E , the V_{noise} is the voltage from the external NOISE unit. The V_{noise} is the frequency dependent signal, filtered from a broadband random noise source by means of a low-pass first-order RC filter at the cut-off frequency of 500 Hz. The root mean square noise voltage is about 10 mV. At $f > 500$ Hz the noise level decreases by a factor of 2 per octave. Thus, it can be considered as $1/f$ noise. For $R_2 = R_3 = 5.6$ k Ω and $E = 10$ V, used in the model circuit, the $V_{on} = 11.5$ V, the $V_{off} = 5.1$ V (assuming $|V_{noise}| \ll E$).

The Schmitt trigger and the electronically controlled switch implement a unit with an S -type current-voltage

characteristic, typically observed in a transistor operating in an avalanche mode [1]–[6]. The subcircuit C2–R4–L–C3–R5 is inserted in the model circuit to imitate the inertial properties of the avalanche transistor. The values of the elements of the circuit in Fig. 6 are: $C_1 = 22$ nF, $R_1 = 24$ k Ω , $L = 5$ mH, $C_2 = C_3 = 10$ nF, $R_4 = R_5 = 510$ Ω .

All variable voltages from the model circuit have been recorded by means of a digital oscilloscope Tektronix TDS2014B (bandwidth 100 MHz). The power spectra have been measured using a low frequency analog spectrum analyzer with the spectral resolution of 10 Hz. The voltage across capacitor C1 saturates to V_0 , which is set slightly less than the ON threshold input voltage of the Schmitt trigger, $V_0 \leq V_{on}$. The trigger waits for a noise signal of *negative* sign to lower, according to (5), the V_{on} so that $V_0 \geq V_{on}$. Then the output of the trigger changes from high to low level and turns on the electronic switch. The latter discharges capacitor C1 and produces the Gaussian shaped output pulse (Fig. 7) with $\dagger \approx 10$ μ s.

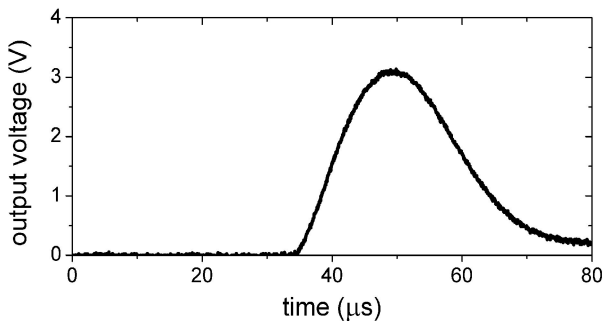


Fig. 7. Individual pulse from the model circuit in Fig. 6.

At higher voltages, $V_0 > V_{on}$ the influence of noise is unimportant and periodic pulses are observed (Fig. 8).

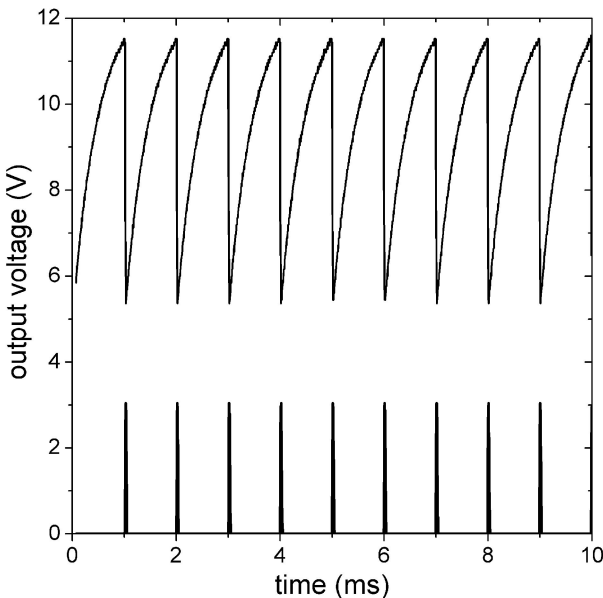


Fig. 8. Periodic pulses from the model circuit; voltage across capacitor C1, taken from the output of the buffer OA (top trace), and output voltage (bottom trace). $V_0 = 12.5$ V.

The random pulse trains, shown in Fig. 9(a) and Fig. 9(b), are observed at supply voltages V_0 very close to the threshold voltage, $V_0 \approx V_{on} = 11.5$ V. Since noise is a random process, the time intervals between the output pulses are also random, providing broadband continuous spectrum

(Fig. 10). The time parameter in (3) of an individual pulse $\dagger = 10$ μ s. Formula (4) gives an estimate, that $ENR(f)$ decreases to the level of 0.8 at $f = 10$ kHz.

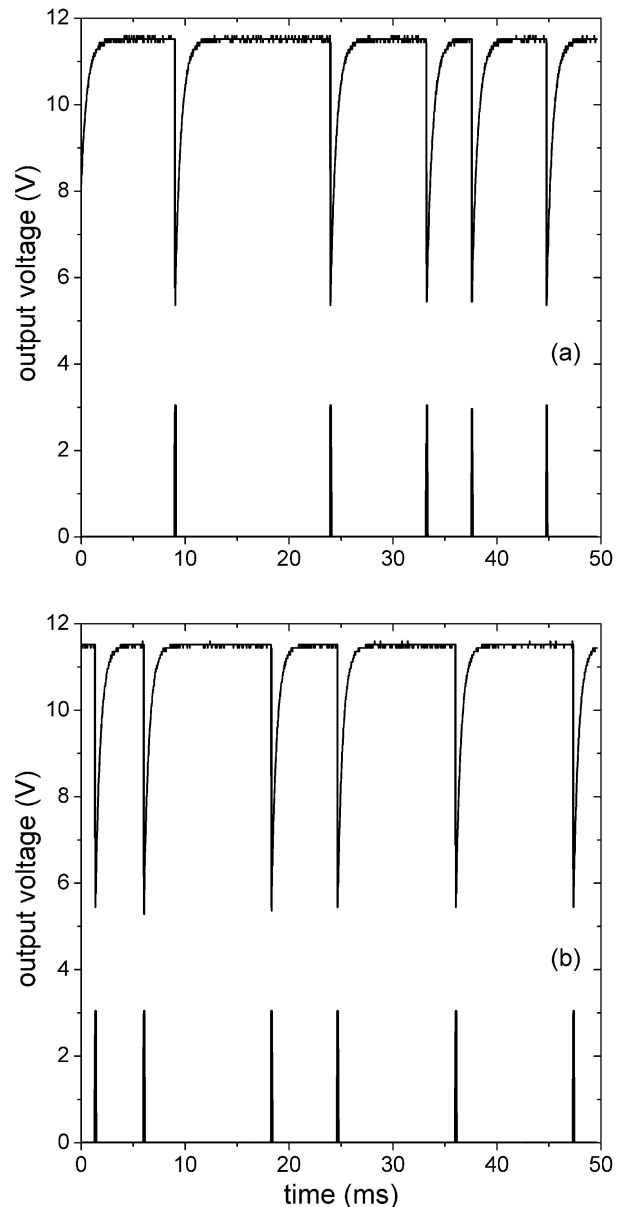


Fig. 9. Two snapshots (a) and (b) of random pulses from the model circuit; voltage across capacitor C1, taken from the output of the buffer OA (top traces) and output voltage (bottom traces). $V_0 = 11.5$ V.

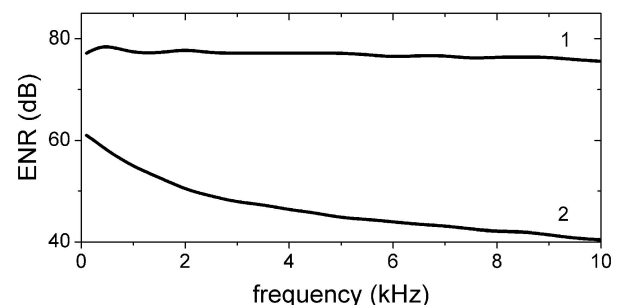


Fig. 10. Power spectra from the model circuit; spectrum of the output signal (1), and spectrum of the noise unit (2). $V_0 = 11.5$ V.

IV. CONCLUSIONS

We have demonstrated, that broadband high spectral density noise with ‘excess noise ratio’ value up to 80 dB, covering the very high (VHF: 30 MHz–300 MHz) and the

ultrahigh (UHF: 300 MHz–1000 MHz) frequency bands, can be generated by an avalanche transistor circuit operating in a random pulse mode. The excess noise ratio is much higher than of many other noise sources and is only by 10 dB lower than of the chaotic waveform oscillators, e.g. the two-stage Colpitts circuit [11], [12]. Whereas, the avalanche transistor random pulse circuit has a big advantage over the chaotic microwave oscillators in respect of the spectrum flatness. Chaotic oscillators generate very high power, up to 90 dB, but have a very serious shortcoming, because of the nonuniformity of the spectra, typically of about 20 dB to 30 dB.

Some recently described chaotic oscillators demonstrate rather uniform spectra [14], [15]. However, their versions, developed so far, operate at low (kilohertz) frequencies only. The avalanche transistor circuit, presented in this paper, exhibits high noise level in the microwave range and good flatness of the spectrum with the nonuniformity of ± 1 dB.

An important observation, made by Dyakonov *et al.* [6], is that not only the transistors, specially designed as the avalanche devices, but many bipolar junction transistors exhibit good performance in the avalanche breakdown region. The silicon bipolar microwave BFG520 type transistor, investigated in this paper, is another example.

To explain the experimental results we have suggested and investigated an analog model, specifically a low frequency electrical circuit, which generates random signals very similar to the pulses, observed experimentally from the avalanche transistor circuit. The time scale in the model circuit, compared to the experimental one, has been changed by a factor of 10^5 : the time parameter τ is increased from 100 ps to 10 μ s. Consequently, the corresponding cut-off frequency, where the spectrum drops to the level of 0.8, is decreased from 1 GHz to 10 kHz. However, the shape of the individual pulse (Fig. 7) and the shape of the power spectrum (Fig. 10, curve 1), obtained from the analog model, agree very well with the experimental plots, shown in Fig. 4 and Fig. 5, respectively.

Similar approach using analog modelling method has been introduced by Namayunas *et al.* [16] and has exhibited very good agreement with the experimental observations. The analog modelling technique (in a sense an analog computer) can be applied to microplasma noise phenomenon [17] to explain the effect of noise enhancing by means of breakdown stimulation with microwave electric field in

avalanche diodes.

REFERENCES

- [1] D. J. Hamilton, J. F. Gibbons, W. Shockley, "Physical principles of avalanche transistor pulse circuits", in *Proc. IRE*, vol. 47, no. 6, 1959, pp. 1102–1108. [Online]. Available: <http://dx.doi.org/10.1109/isscc.1959.1157029>
- [2] W. M. Henebry, "Avalanche transistor circuits", *Rev. Sci. Instr.*, vol. 32, no. 11, pp. 1198–1203, 1961.
- [3] P. Spirito, G. F. Vitale, "An analysis of the dynamic behavior of switching circuits using avalanche transistors", *IEEE J. Solid-State Circuit*, vol. SC-7, no. 4, pp. 315–320, 1972.
- [4] V. P. Dyakonov, "Nanosecond rectangular pulse generators using avalanche and MIS transistors", *Instrum. Exp. Tech.*, vol. 23, no. 4, pp. 913–914, 1980.
- [5] V. P. Dyakonov, "Nanosecond pulse generators based on avalanche and MIS transistors", *Instrum. Exp. Tech.*, vol. 24, no. 1, pp. 131–132, 1981.
- [6] V. P. Dyakonov, T. A. Vasilkova, Yu. A. Ermachkova, "Measurement of the pulse parameters of silicon transistors under avalanche operating conditions", *Meas. Tech.*, vol. 50, no. 7, pp. 770–774, 2007.
- [7] G. Duan, S. Vainshtein, J. Kostamovaara, "Three-dimensional peculiarities in an avalanche transistor provide a broadened range of amplitudes and durations in the generated pulses", *Appl. Phys. Lett.*, vol. 101, no. 17, pp. 173–506, 2012. [Online]. Available: <http://dx.doi.org/10.1063/1.4764114>
- [8] TD/TN series microwave noise tubes and noise sources. [Online]. Available: <http://www.highenergydevices.com>
- [9] Calibrated sources: NC3000 coaxial. [Online]. Available: <http://noisecom.com>
- [10] A. Namajunas, A. Tamasevicius, G. Mykolaitis, S. Bumeliene, J. Pozela, "Microplasma noise stimulated by microwave electric field", *Acta Phys. Polon. A*, vol. 107, no. 2, pp. 369–372, 2005.
- [11] A. Tamasevicius, G. Mykolaitis, S. Bumeliene, A. Cenys, A. N. Anagnostopolous, E. Lindberg, "Two-stage chaotic Colpitts oscillator", *Electron. Lett.*, vol. 37, no. 9, pp. 549–551, 2001.
- [12] S. Bumeliene, A. Tamasevicius, G. Mykolaitis, A. Baziliauskas, E. Lindberg, "Numerical investigation and experimental demonstration of chaos from two-stage chaotic Colpitts oscillator in the ultrahigh frequency range", *Nonlinear Dyn.*, vol. 44, no. 1–4, pp. 167–172, 2006.
- [13] P. Horowitz, W. Hill, *The Art of Electronics*. Cambridge, New York: Cambridge University Press, 1993, pp. 231–232.
- [14] A. Tamasevicius, S. Bumeliene, R. Kirvaitis, G. Mykolaitis, E. Tamaseviciute, E. Lindberg, "Autonomous Duffing–Holmes type chaotic oscillator", *Elektronika ir Elektrotechnika*, no. 5, pp. 43–46, 2009.
- [15] E. Tamaseviciute, G. Mykolaitis, A. Tamasevicius, "Autonomous Silva–Young chaotic oscillator with flat power spectrum", *Elektronika ir Elektrotechnika*, no. 9, pp. 109–112, 2011.
- [16] A. M. Namayunas, Yu. K. Pozhela, A. V. Tamashyavichyus, "Microplasma instability in alternating electric field", *Sov. Phys. Dokl.*, vol. 34, pp. 630–632, 1989.
- [17] J. Pozhela, *Plasma and Current Instabilities in Semiconductors*, Oxford, New York: Pergamon, 1981.